

Transition to the Matter Dominated Era, Recombination, and the Beginning of Chemistry

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Outline

1. Matter Domination
2. Recombination: Equilibrium Model
3. Recombination: Kinetics
4. The Post-Recombination Universe

The End of the Radiation Era

- Radiation energy density (photons + neutrinos):

$$\rho_r = 7.8 \times 10^{-34} (1 + z)^4 \text{ g/cm}^3$$

- Compare to matter energy density:

$$\rho_m = 1.88 \times 10^{-29} \Omega_m h^2 (1 + z)^3 \text{ g/cm}^3$$

- The matter becomes dominant at:

$$1 + z_{\text{eq}} = 24100 \Omega_m h^2 \approx 3100$$

Recombination

- Recall Saha equation (the original).

$$\frac{n_e n_p}{n(\text{H I})} = \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{-I/kT}$$

- $I = 13.6 \text{ eV}$, $I/k = 158000 \text{ K}$.
- Define $x_e = n_e/n_H$ (all H nuclei). Then:

$$\frac{x_e x_p}{x(\text{H I})} = \frac{1}{n_H} \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{-I/kT}$$

$$n_H = 1.0 \times 10^4 \frac{\Omega_b h^2}{0.023} \frac{1-Y}{0.76} T_4^3 \text{ cm}^{-3}$$

Recombination, Part 2

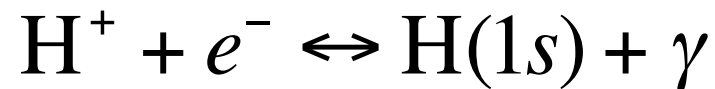
- In thermal equilibrium, the balance of ionized (p+e) and neutral (HI) hydrogen is:

$$\frac{x_e x_p}{x(\text{H I})} = 2.5 \times 10^{17} \frac{0.023}{\Omega_b h^2} \frac{0.76}{1-Y} T_4^{-3/2} e^{-15.8/T_4}$$

- Universe recombines when exponential factor outweighs the 10^{17} .
- 50% ionization expected @ 3740 K, $z=1370$.

Reaction Kinetics

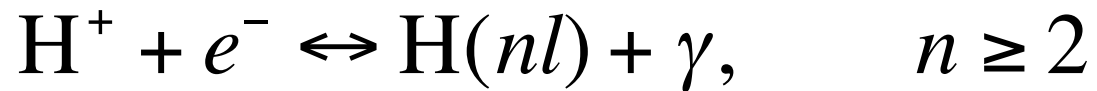
- In the real Universe, recombination is not an equilibrium process. Must identify reactions that lead to recombination.
- First try:



- Doesn't work. Universe is optically thick to $E > 13.6\text{eV}$ photons if $x_{\text{HI}} > 10^{-8}$.

Recombination via Excited States

- First recombine to excited states – Universe is optically thick to these photons

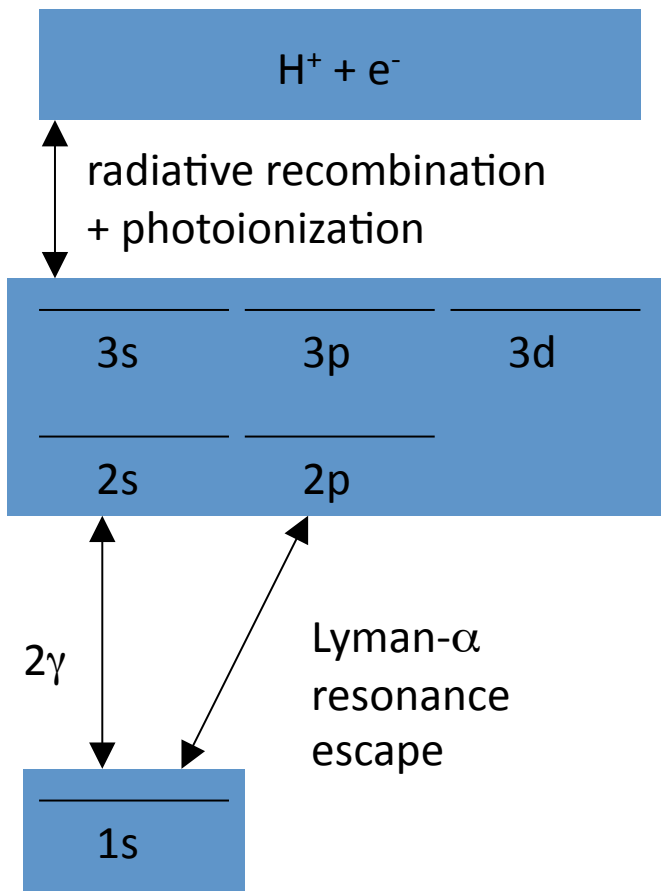


- Decay to ground state:

$\text{H}(2p) \leftrightarrow \text{H}(1s) + \gamma_{\text{Ly}\alpha}$, photon redshifts out of line

$\text{H}(2s) \leftrightarrow \text{H}(1s) + \gamma + \gamma$

Recombination via Excited States



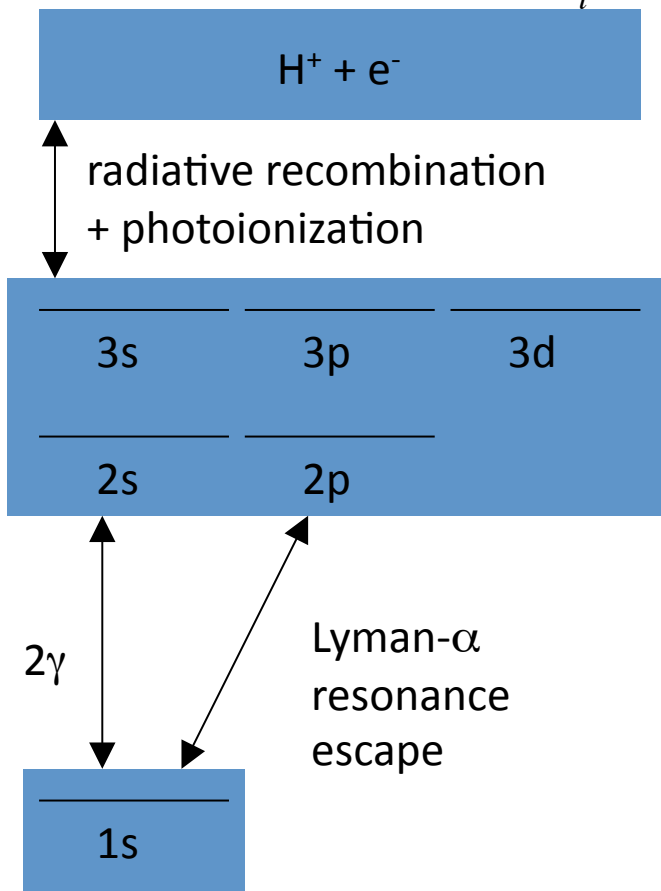
- Effective recombination rate is recombination coefficient to excited states times branching fraction to ground state.
- Must include ionization rate to achieve balance in the case of Saha equilibrium.

$$\frac{\# \text{ rec}}{\Delta V \Delta t} = \frac{2\Lambda + 6A_{Ly\alpha}P_{esc}}{2\Lambda + 6A_{Ly\alpha}P_{esc} + \sum_i g_i e^{-(E_i - E_2)/kT} \beta_i} \alpha_B n_e n_p$$

$$\alpha_B = \sum_{nl, n \geq 2} \alpha_{nl}$$

Recombination via Excited States

$$\frac{dx_{HI}}{dt} = \frac{2\Lambda + 6A_{Ly\alpha}P_{esc}}{2\Lambda + 6A_{Ly\alpha}P_{esc} + \sum_i g_i e^{-(E_i - E_2)/kT} \beta_i} \alpha_B \left[x_e x_p n_H - \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{-I/kT} x_{HI} \right]$$



Λ = 2-photon decay rate from 2s = 8.2 s^{-1}

P_{esc} = escape probability from Lyman- α line

$A_{Ly\alpha}$ = Lyman- α decay rate = $6 \times 10^8 \text{ s}^{-1}$.

α_e = recombination rate to excited states

g_i = degeneracy of level $l = 2l_i + 1$

β_i = photoionization rate from level l

I = ionization energy

Escape probability

- Total optical depth for a photon redshifting through Lyman- α line:

$$\sigma(\nu) = \frac{3\lambda_{Ly\alpha}^2 A_{Ly\alpha}}{8\pi} \delta(\nu - \nu_{Ly\alpha})$$

$$\tau = \int n_{HI} \sigma c dt = n_{HI} c \int \sigma(\nu) \frac{d\nu}{H\nu} = \frac{3\lambda_{Ly\alpha}^3 A_{Ly\alpha} n_H x_{HI}}{8\pi H}$$

- Usual (Sobolev) approximation for P_{esc} :
 - Emission and absorption line profiles are the same
 - So to escape a photon has to traverse an optical depth τ_{esc} uniformly distributed between 0 and τ .

$$P_{esc} = \langle e^{-\tau_{esc}} \rangle = \frac{1}{\tau} \int_0^\tau e^{-\tau_{esc}} d\tau_{esc} = \frac{1 - e^{-\tau}}{\tau}$$

Simplifications to Recombination Equation

- Escape Probability: practical case: $\tau \gg 1$, so

$$P_{esc} \approx \frac{1}{\tau} = \frac{8\pi H}{3\lambda_{Ly\alpha}^3 A_{Ly\alpha} n_H x_{HI}}$$

- Ionization rate: from detailed balance: (recall energy of 2nd level is $-I/4$)

$$\sum_i g_i e^{-(E_i - E_2)/kT} \beta_i = 2\alpha_B \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{I/4kT}$$

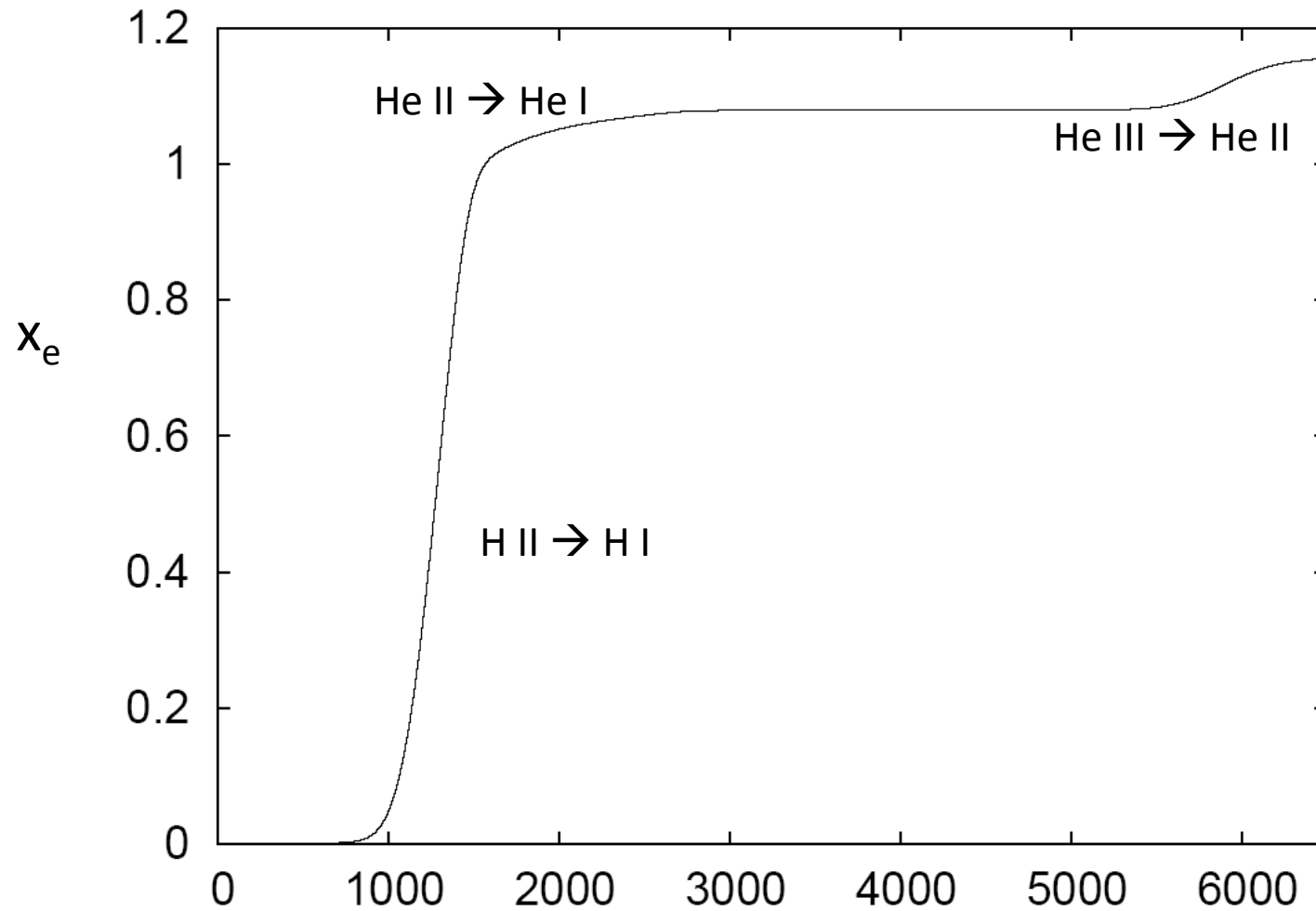
- This leads to the Peebles (1968) or Zel'dovich et al (1968) equation.

The Peebles/Zel'dovich et al ODE

$$\frac{dx_{HI}}{dt} = \frac{\Lambda + \frac{8\pi H}{\lambda_{Ly\alpha}^3 n_H x_{HI}}}{\Lambda + \frac{8\pi H}{\lambda_{Ly\alpha}^3 n_H x_{HI}} + \alpha_B \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{I/4kT}} \alpha_B \left[x_e x_p n_H - \left(\frac{m_e kT}{2\pi\hbar^2} \right)^{3/2} e^{-I/kT} x_{HI} \right]$$

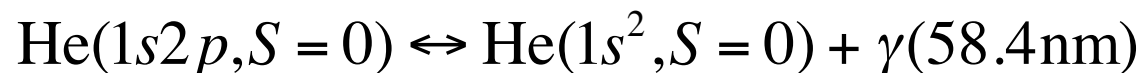
- Two limiting cases:
 - “Early”: Decay from excited states is the bottleneck ($z > 850$)
 - “Late”: Recombination to excited state is the bottleneck, also ionization negligible ($z < 850$).

Solution

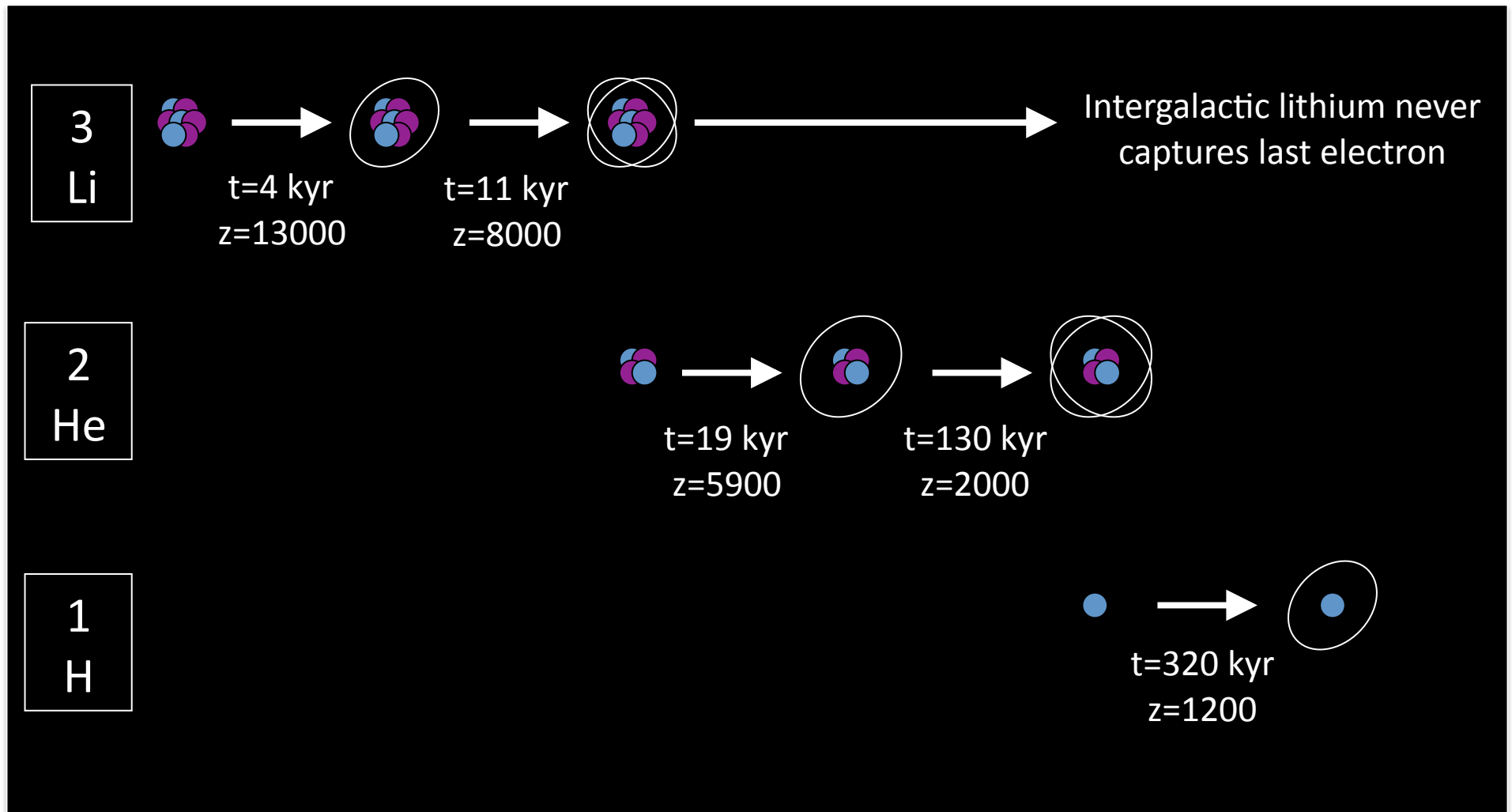


Helium recombination

- He:H=0.08 by number, so for fully ionized gas $x_e=1.16$.
- Helium recombination similar to hydrogen:
 - He III \rightarrow He II @ $z \sim 5800$ ($I = 54.4$ eV).
 - He II \rightarrow He I @ $z \sim 2000$ ($I = 24.6$ eV) has additional pathways to ground state since at $z < 2200$ there is a small amount of H I present.



Recombination Timeline



Implications for CMB

- Optical depth to Thomson scattering per e-fold of expansion is:

$$\frac{d\tau}{d\ln a} = \frac{n_H x_e \sigma_T c}{H} \approx 3.8 \times 10^{-3} (1+z)^{3/2} x_e$$

- At recombination era, $z \sim 10^3$ so this is $\gg 1$ when $x_e \sim 1$. Recombination makes the universe transparent.
- Integrated optical depth is unity at $z \sim 1100$. (“Surface of last scattering.”)

After Recombination

- Tail end of recombination is second-order reaction. Once ionizations become insignificant:

$$\frac{dx_e}{dt} = -\alpha_B n_H x_e^2$$

$$x_e^{-1} = \text{const} + \int \alpha_B n_H dt$$

- Roughly, $\alpha_B \sim 10^{-12} \text{ cm}^3/\text{s}$, $n_H \sim 100 \text{ cm}^{-3}$, $t \sim 10^{13} \text{ s}$, so expect recombination to end with $x_e \sim 10^{-3}$. Detailed calculations give 2×10^{-4} (e.g. Scott & Moss 2009).

Thermal Decoupling

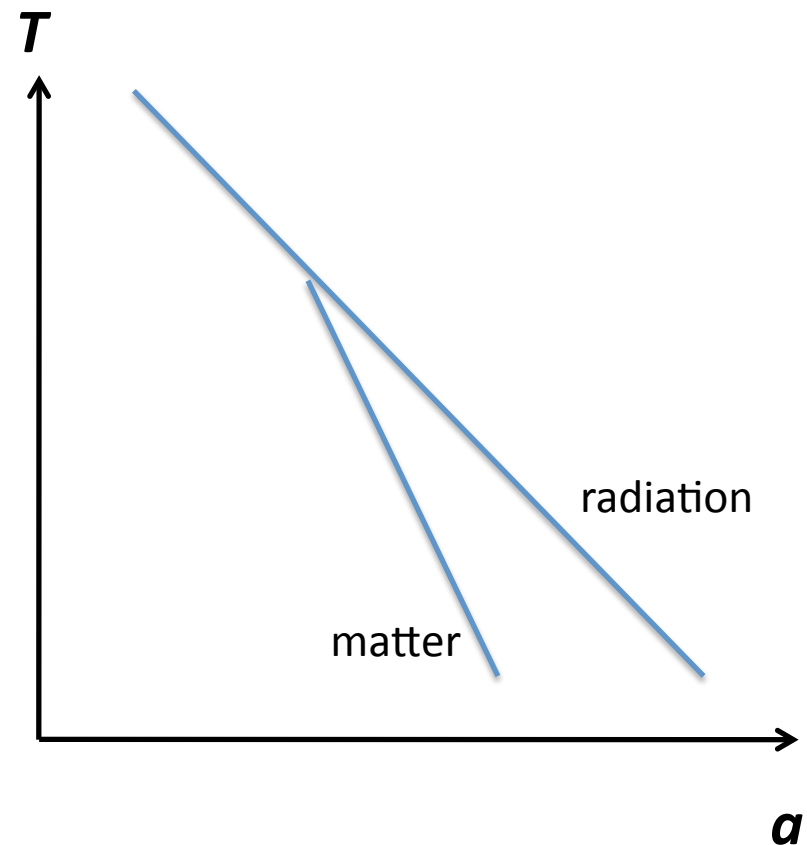
- Recall the timescale to equilibrate the baryon temperature with the CMB radiation:

$$t_b \sim \frac{m_e c^2 \langle E_\gamma \rangle^{-1}}{n_\gamma \sigma_T c x_e} = \frac{m_e c}{\sigma_T u_\gamma x_e} = 2 \times 10^{13} T_3^{-4} \text{ sec}$$

- Now with x_e factor since only free electrons contribute to Compton equilibrium.
- Compare to age of Universe which scales as $T^{-3/2}$ in matter era.
- There comes a time when the baryonic gas is no longer coupled to the CMB temperature. $z=126$, $T=346$ K (Scott & Moss 2009).

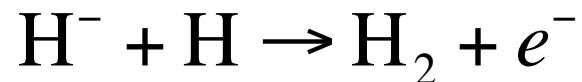
Thermal Decoupling, Part 2

- Gas expands adiabatically. $T \sim \rho^{2/3}$ for monatomic gas, so the matter temperature $T_m \sim a^{-2}$.
- The matter cools below the CMB temperature, e.g. reaches 10 K at $z=20$ if there is no heat injection.



Molecules?

- At $z < 400$, formation of H_2 is thermodynamically favored.
- Reactions are slow however:



- Pregalactic gas $H_2:H \sim 6 \times 10^{-7}$ (Hirata & Padmanabhan 2006).

Speed of Sound

- General: $c_s^2 = \left. \frac{\partial p}{\partial \rho} \right|_S$
- Before recombination:
 - Photons & baryons are tightly coupled. If compressed by $\delta \ln V$, we get:

$$\delta p = \delta p_r = \frac{1}{3} \delta u_r = -\frac{4}{9} u_r \delta \ln V$$

$$\delta \rho = \frac{\delta u_r}{c^2} + \delta \rho_b = -\left(\frac{4u_r}{3c^2} + \rho_b \right) \delta \ln V$$

$$c_s^2 = \frac{\frac{4}{9} u_r}{\frac{4u_r}{3c^2} + \rho_b} = \frac{c^2}{3\left(1 + \frac{3\rho_b c^2}{4u_r}\right)} \approx \frac{c^2}{3}$$

Speed of Sound

- After recombination:
monatomic gas, $p \sim \rho^{5/3}$:

$$c_s = \sqrt{\frac{5kT}{3m_H}} = 3.7 \times 10^5 T_3^{1/2} \text{ cm/s}$$

- Break in sound speed at $z=126$, $c_s=2 \times 10^5 \text{ cm/s}$ due to thermal evolution.
- Sound speed in DM is always zero.

